on the basis of energy calculations for the nonprotonated molecule (Chachaty, Pappalardo \& Scarlata, 1976). On the other hand, theoretical calculations (Galasso, De Alti \& Bigotto, 1971) and molar Kerr constant studies (Aroney, Le Fevre \& Saxby, 1963; Le Fevre \& Saxby, 1966) predict for the diphenyl sulphide molecule a non-planar structure with angles of $c a 42-$ $48^{\circ}$ between the phenyl ring planes because of the steric contact between ortho H atoms. It is therefore reasonable that in this case the presence of a single H atom in the ortho position has too little steric influence to alter significantly the planar structure of the molecule. In addition, the contact distance of $2.63 \AA$ between the two N atoms strongly suggests the existence of an intramolecular $\mathrm{N}-\mathrm{H} \cdots \mathrm{N}$ hydrogen bond. The $\mathrm{C}-\mathrm{S}-\mathrm{C}$ angle of $108(1)^{\circ}$ shows that the S atoms are $s p^{3}$-hybridized. The mean value of the two $\mathrm{C}-\mathrm{S}$ bond lengths is $1.755 \AA$, which is appreciably shorter than the paraffinic distance of $1.81 \AA$, but is in agreement with the $\mathrm{C}\left(s p^{2}\right)-\mathrm{S}$ bond with two atoms bonded to the S atoms, e.g. $1.77 \AA$ in phenothiazine (Bell, Blount, Briscoe \& Freeman, 1968), and $1.765 \AA$ in methixene (Chu, 1972). The relative shortness of the $\mathrm{S}-\mathrm{C}$ bonds could indicate involvement of the sulphur $d$ orbitals with the pyridine $\pi$ systems and via hyperconjugation they also contribute to the cation's planarity. Bond distances and angles in the pyridine rings have normal values. As shown in the [001] projection of the unit cell (Fig. 2) the entire structure comprises alternate cationic and anionic units developing along [100]. The thickness
of the $\left[\mathrm{UO}_{2} \mathrm{Cl}_{4}\right]^{2-}$ groups and of the protonated ligands is roughly 3.3 and $5 \AA$, respectively, in the $\mathbf{b}$ direction.

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# (2-Oxopropyl)mercury Bromide 

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#### Abstract

C}_{3} \mathrm{H}_{5} \mathrm{OHgBr}\), triclinic, $P \overline{1}, a=6.752$ (1), $b=$ 5.733(1), $c=7.979$ (1) $\AA, \quad \alpha=83.14$ (1), $\beta=$ 95.53 (1), $\gamma=101.47(1)^{\circ}, Z=2, D_{o}=3.70(5), D_{c}=$ $3.741 \mathrm{~g} \mathrm{~cm}^{-3}$. The structure contains two short intermolecular $\mathrm{Hg} \cdots \mathrm{O}(\mathrm{CO})$ contacts [2.81 (2), 2.82 (2) $\AA$ ] which are consistent with the relatively low carbonyl stretching frequency observed ( $1645 \mathrm{~cm}^{-1}$ ), assuming that formation of weak intermolecular $\mathrm{Hg} \cdots \mathrm{O}$ bonds


also weakens the carbonyl bond. The $\mathrm{Hg}-\mathrm{Br}$ and $\mathrm{Hg}-\mathrm{C}$ distances are 2.441 (5) and 2.08 (4) $\AA$, respectively, while the $\mathrm{C}-\mathrm{Hg}-\mathrm{Br}$ angle is $176(1)^{\circ}$.

Introduction. Milky-white, irregularly shaped crystals of (2-oxopropyl)mercury bromide, obtained from acetone-petroleum ether, were provided by Mr John Feeney (1976). A large crystal was air blown around a

Carborundum ring to yield a sphere of diameter 0.16 mm which was glued to the end of a glass rod and centered on a Syntex $P 2_{1}$ computer-controlled diffractometer. Diffractometer examination of the reciprocal lattice revealed no systematic absences. Using a $\theta-2 \theta$ scan and graphite-monochromated Mo $K \alpha$ radiation ( $\lambda=0.71069 \AA$ ), intensity data were collected in the range $4<2 \theta<60^{\circ}$. Three standard reflections, measured at 47 -reflection intervals, remained constant to $\pm 5 \%$ throughout the data collection period. Of the 1746 unique reflections measured, 771 with $F^{2} \geq 3 \sigma\left(F^{2}\right)$ (counting statistics) were used in the structure solution and refinement. Spherical absorption corrections (International Tables for X-ray Crystallography, 1972) were applied to all data with $\mu$ for Mo $K \alpha=331 \mathrm{~cm}^{-1}$.

The structure was solved by the heavy-atom method and refined using full-matrix least-squares techniques; the computer programs used have been described previously (Potenza, Giordano, Mastropaolo \& Efraty, 1974). Approximate coordinates for the Hg and Br atoms were determined readily from a normal sharpened Patterson map. The distribution of peaks in this map arising from vectors between the heavy atoms suggested that the structure was centrosymmetric. Oxygen and C atom coordinates were determined from an electron-density map based on phases from the heavy atoms. No attempt was made to locate H atoms. Neutral-atom scattering factors were obtained from Cromer \& Waber (1965); both real and imaginary components of the anomalous dispersion factors (International Tables for X-ray Crystallography, 1962) were applied to Hg and Br .

Refinement was based on $F$ and the function minimized was $R_{w F}=\left[\sum w\left(\left|F_{o}\right|-\left|F_{c}\right|\right)^{2} / \sum w F_{o}^{2}\right]^{1 / 2}$ where $w=1 / \sigma^{2}$. For the final refinement cycles, weights were determined by an analysis of variance to make $|\Delta F| / \sigma$ independent of $\left|F_{o}\right|$. This led to the following assignments: $\sigma\left(F_{o}\right)=-0.058\left|F_{o}\right|+11.60$ for $\left|F_{o}\right|<84.7$, $\sigma\left(F_{o}\right)=0.038\left|F_{o}\right|+3.07$ for $84.7<\left|F_{o}\right|<131.7$ and $\sigma\left(F_{o}\right)=0.138\left|F_{o}\right|-11.02$ for $\left|F_{o}\right|>131.7$. Utilizing anisotropic thermal parameters for all atoms, the structure converged with $R_{F}=\sum| | F_{o}\left|-\left|F_{c}\right|\right| \sum\left|F_{o}\right|=$ 0.090 and $R_{w F}=0.101$.* For the final refinement cycle, all parameter changes were less than $0.55 \sigma$, where $\sigma$ is the e.s.d. obtained from the inverse matrix. A final difference map showed no extraneous peaks. Atomic parameters are shown in Table 1 while a view of the molecule, showing bond distances and angles, is given in Fig. 1. A stereoview of the molecular packing is given in Fig. 2.

[^0]Excluding atom $\mathrm{C}(2)$, r.m.s. thermal displacements along the principal axes of the ellipsoids are reasonable and vary from 0.17 to $0.36 \AA$. $\mathrm{C}(2)$ shows an unusually small value of $0.06 \AA$ along one axis. The reasons for this are not entirely clear and may be related to the relatively large value of the linear absorption coefficient of the compound (uncompensated absorption). We do note, however, that the ellipsoid for $\mathrm{C}(2)$ does 'point' in a reasonable direction, i.e. approximately normal to the various bond axes.

Discussion. A number of $\alpha$-substituted ketone and aldehyde mercurials have been synthesized and characterized by IR spectroscopy (Allmann, Flatau \& Musso, 1972; Feeney, 1976). In each case, the carbonyl stretching frequency was significantly lower than that found for the free aldehyde or ketone. Examination of the IR spectrum of ( 2 -oxopropyl)mercury bromide revealed a band centered at $1645 \mathrm{~cm}^{-1}(\mathrm{KBr})$ which, although low in frequency, was assigned as a carbonyl stretching mode. The present structural study was undertaken to characterize the nature of the interaction leading to this low frequency.

Structural parameters of the molecule are in general agreement with accepted values. All C and O atoms are planar to within $\pm 0.02 \AA$ while the bond angles with $C(2)$ as vertex range from $118(3)$ to $121(3)^{\circ}$,


Fig. 1. (2-oxopropyl)mercury bromide showing the atom numbering scheme and bond distances. Bond angles are as follows: $\mathrm{Br}-\mathrm{Hg}-\mathrm{C}(1), 176$ (1); $\mathrm{Hg}-\mathrm{C}(1)-\mathrm{C}(2), 108$ (2); $\mathrm{O}-$ $\mathrm{C}(2)-\mathrm{C}(1), 118(3) ; \mathrm{O}-\mathrm{C}(2)-\mathrm{C}(3), 121$ (3); C(1)-C(2)-C(3), $120(3)^{\circ}$.

Table 1. Final atomic parameters
Estimated standard deviations, obtained from the least-squares refinement, are given in parentheses. Coordinate values are $\times 10^{3}$.

|  | $x$ | $y$ | $z$ |
| :--- | :---: | :---: | :---: |
| Hg | $43 \cdot 4(2)$ | $861 \cdot 2(2)$ | $-218 \cdot 1(2)$ |
| Br | $-223 \cdot 7(6)$ | $1056 \cdot 2(8)$ | $-365 \cdot 3(5)$ |
| O | $194(4)$ | $283(4)$ | $-60(3)$ |
| $\mathrm{C}(1)$ | $279(6)$ | $717(8)$ | $-83(5)$ |
| $\mathrm{C}(2)$ | $278(4)$ | $469(4)$ | $-146(4)$ |
| $\mathrm{C}(3)$ | $364(8)$ | $456(7)$ | $-306(6)$ |



Fig. 2. Stereoscopic packing diagram of (2-oxopropyl)mercury bromide showing the coordination geometry of Hg. Weak intermolecular $\mathrm{Hg} \cdots \mathrm{O}$ bonds are indicated by light lines.
suggesting $s p^{2}$ hybridization at that atom. The $\mathrm{Br}-\mathrm{Hg}-$ $\mathrm{C}(1)$ angle $\left[176(1)^{\circ}\right]$ is equivalent to the $\mathrm{C}-\mathrm{Hg}-\mathrm{C}$ angle $\left[175.3(3)^{\circ}\right]$ found for $\mathrm{C}\left(\mathrm{HgOCOCF}_{3}\right)_{4}$ (Grdenić, Kamenar, Korpar-Colig, Sikirica \& Jovanovski, 1974) and approximates the $180^{\circ}$ angle expected for free organo-mercurials. The $\mathrm{Hg}-\mathrm{Br}$ distance $[2.441$ (5) $\AA$ ] lies between those of 2.433 (5) and 2.473 (12) $\AA$ reported for $\mathrm{HgBr}_{2} \cdot 2 \mathrm{C}_{4} \mathrm{H}_{8} \mathrm{O}_{2}$ and $\mathrm{HgBr} . \mathrm{C}_{4} \mathrm{H}_{8} \mathrm{O}$, respectively (Frey \& Monier, 1971), while the $\mathrm{Hg}-\mathrm{C}(1)$ distance lies at the short end of the range $(2.08-2 \cdot 15 \AA)$ found with several organomercury compounds (Grdenić, Kamenar \& Nagl, 1977).

The tendency of Hg to coordinate additional atoms weakly in the solid state has been noted by Grdenic (1965). In the present structure, molecules are oriented such that the shortest intermolecular contacts involving Hg occur with O atoms on molecules related by a cell translation along $b\left[\mathrm{Hg} \cdots \mathrm{O}^{\prime}, 2 \cdot 82\right.$ (2) $\AA$ ] and by molecules related by inversion followed by a cell translation along $b\left[\mathrm{Hg} \cdots \mathrm{O}^{\prime \prime}, 2.81\right.$ (2) $\AA$ ]. These distances are slightly shorter than the sum of the van der Waals radius of O and the metallic radius of divalent $\mathrm{Hg}(2.90$ $\AA$ ). They are approximately the same as that separating different $\mathrm{Hg}-\mathrm{O}$ chains in the hexagonal form of HgO [2.79 (2) $\AA$, Aurivillius \& Carlsson (1958)] and significantly longer than the $\mathrm{Hg} \cdots \mathrm{O}$ distances of 2.70 (2) and 2.56 (3) $\AA$ reported, respectively, for bis(dipivaloylmethyl)mercury (Allmann et al., 1972) and $\mathrm{C}_{7} \mathrm{H}_{7} \mathrm{OHg} . \mathrm{C}_{7} \mathrm{H}_{6} \mathrm{O}_{2} \mathrm{~N}$ (Kobayashi, Iitaka \& Kido, 1970). With HgO , this short distance is thought to imply a weak bonding interaction (Cotton \& Wilkinson, 1972). The short $\mathrm{Hg} \cdots \mathrm{O}$ contacts in the title complex are consistent with the low carbonyl stretching frequency observed if it is assumed that they imply weak $\mathrm{Hg} \cdots \mathrm{O}$ bonding which also weakens the carbonyl bond. No other contacts involving Hg are short
enough to imply electronic interaction stronger than that arising from van der Waals forces. Assuming $\mathrm{O}^{\prime}$ and $\mathrm{O}^{\prime \prime}$ to be weakly coordinated to Hg , the coordination geometry of Hg is that of an octahedron with two cis sites vacant. The angles between cis ligands with Hg as vertex range from 85 to $90^{\circ}$.

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[^0]:    * Lists of structure factors and anisotropic thermal parameters have been deposited with the British Library Lending Division as Supplementary Publication No. SUP 33542 ( 7 pp .). Copies may be obtained through The Executive Secretary, International Union of Crystallography, 5 Abbey Square, Chester CH1 2HU, England.

